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## ภาคผนวก

ผลงานตีพิมพ์ในวารสารวิชาการจากโครงการวิจัย

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# Spatial Profiles of Solar Storm Particles during Finite-Time Shock Acceleration

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## ABSTRACT

We apply a model of finite time shock acceleration for energetic storm particle events at interplanetary shocks to describe the spatial profiles of energetic ions during events in September 1999 and October 2000. The finite time shock acceleration model is a development of the diffusive shock acceleration mechanism, incorporating rigidity-dependent scattering due to magnetic turbulence in astrophysical shock phenomena. A system of coupled ordinary differential equations is solved to describe the time evolution of the particle density at tens or hundreds of energy values. A simultaneous fit to the spectra of carbon, oxygen, and iron ions is used to determine a key parameter, the scattering mean free path  $\lambda$ , as a function of particle rigidity. Next, we model the spatial profiles by solving the parabolic partial differential equation for particle diffusion, either analytically (for a steady state) or numerically (for time evolution). We aim to compare the results with measured spatial profiles of solar storm particles from Ultra-Low-Energy Isotope Spectrometer (ULEIS) on board the Advanced Composition Explorer (ACE) spacecraft.

**Keywords:** Shock acceleration, interplanetary shock, energetic storm particle event

## INTRODUCTION

The Earth is impacted by energetic particles from the Sun and Galactic cosmic rays, including particles from various sources in our galaxy. The most energetic particles in the solar system are released from the Sun. The energetic particles from the Sun consist of protons, electrons and heavy ions with energy ranging from a few tens of keV to GeV (the fastest particles can reach speed up to 80% of the speed of light). They are of particular interest and importance because they can endanger life in outer space (especially particles above 40 MeV). Solar energetic particles (SEPs) can originate from two processes: energization at a solar flare site or by shock waves associated with coronal mass ejection (CMEs).

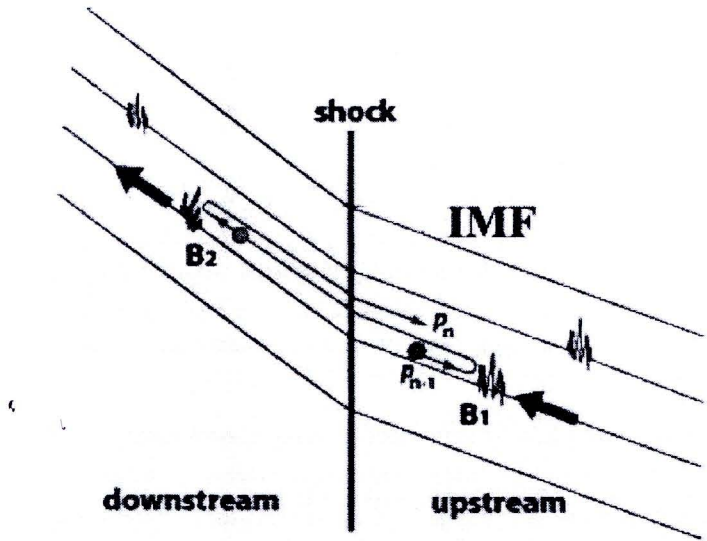
Particle acceleration occurs as a consequence of transient releases of energy in association with solar flares or coronal mass ejections. When a shock wave develops, it expands and propagates through the interplanetary medium. If the magnetohydrodynamic strength of the shock is high enough, it is able to accelerate particles which propagate along the interplanetary magnetic field (IMF) lines (see figure 1).

The observed spectra of energetic particles from the Sun can indicate enhanced intensity in association with the passage of interplanetary shocks. These are known as solar storm particles or energetic storm particles (ESP) [1]. These typically exhibit spectral rollovers at 0.1 to 10 MeV nucleon<sup>-1</sup> [2,3]. The physical origin of such rollovers is the finite time available for shock acceleration.

We apply a model of finite time shock acceleration for energetic storm particle events at interplanetary shocks to describe the spatial profiles of energetic ions during events in solar cycle 23 (during September 1999 and October 2000). The finite time shock acceleration model is a development of the diffusive shock acceleration mechanism, incorporating rigidity-dependent scattering due to magnetic turbulence in astrophysical shock phenomena [4]. A



system of coupled ordinary differential equations is solved to describe the time evolution of the particle density at tens or hundreds of energy values. A simultaneous fit to the spectra of carbon, oxygen, and iron ions is used to determine a key parameter, the scattering mean free path ( $\lambda$ ), as a function of particle rigidity.



**Figure 1.** Acceleration of a particle propagating along the interplanetary magnetic field (IMF) lines crossing the shock (discontinuity of direction of magnetic field,  $B$ ).

### COMPUTATIONAL METHOD

The finite-time shock acceleration (FTSA) model [5] is based on the probabilistic model. The key quantities are the number of acceleration events,  $n$ , the rate of acceleration events,  $r$ , the rate of escape  $\varepsilon$ , and the duration of the shock acceleration process,  $t$ . The mathematical model of the finite-time shock acceleration is a system of coupled ordinary differential equations, which is solved to describe the time evolution of the particle density,  $N(t)$ :

$$\frac{dN_n(t)}{dt} = I_n - (r_n + \varepsilon_n)N_n(t) + r_{n-1}N_{n-1}(t) \quad (1)$$

where  $N_n(t)$  is the particle density,  $r_n$  is the rate of acceleration,  $\varepsilon_n$  is the rate of escape, and  $I_n$  is the rate of particle inflow to the shock (the subscript  $n$  is the number of acceleration events).

We model the spatial profiles by solving the parabolic partial differential equation for particle diffusion, either analytically (for a steady state) or numerically (for time evolution). Consider the diffusion convection equation (equation 2) for describing the particle density,  $F$  at the shock of interplanetary space ( $z = 0$ ).

$$\frac{\partial F}{\partial t} = \kappa \frac{\partial^2 F}{\partial z^2} - U \frac{\partial F}{\partial z} \quad (2)$$

where  $F(z,t)$  is the particle density at position  $z$  at time  $t$ ,  $U$  is the fluid speed, and  $\kappa$  is the spatial diffusion coefficient.

We set a boundary condition at the shock of  $F(z=0) = N_0 \exp(t/\tau)$  and at positions far from the shock of  $F(z \rightarrow \infty) = F(z \rightarrow -\infty) = N_0$ , where  $N_0$  is the seed particle density and  $\tau$  is the exponential time scale. For the initial condition we set  $F(z, t=0) = N_0$ .

We solve equation (2) with finite differencing by the Crank-Nicolson Method. We write it in discrete form as

$$\begin{aligned} \frac{F_i^{n+1} - F_i^n}{\Delta t} = & \kappa \frac{1}{\Delta z^2} \left[ \frac{(F_{i+1}^{n+1} - 2F_i^{n+1} + F_{i-1}^{n+1}) + (F_{i+1}^n - 2F_i^n + F_{i-1}^n)}{2} \right] \\ & - U \frac{1}{\Delta z} \left[ \frac{(F_{i+1}^{n+1} - F_{i-1}^{n+1}) + (F_{i+1}^n - F_{i-1}^n)}{4} \right]. \end{aligned} \quad (3)$$

Then

$$\begin{aligned} F_i^{n+1} = & F_i^n + \kappa \frac{\Delta t}{\Delta z^2} \left[ \frac{(F_{i+1}^{n+1} - 2F_i^{n+1} + F_{i-1}^{n+1}) + (F_{i+1}^n - 2F_i^n + F_{i-1}^n)}{2} \right] \\ & - U \frac{\Delta t}{\Delta z} \left[ \frac{(F_{i+1}^{n+1} - F_{i-1}^{n+1}) + (F_{i+1}^n - F_{i-1}^n)}{4} \right] \end{aligned} \quad (4)$$

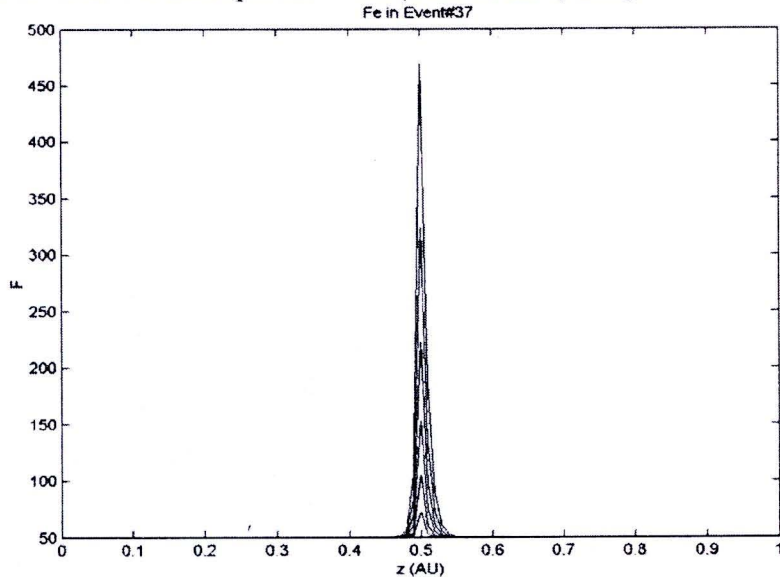
and we get

$$\begin{aligned} F_i^{n+1} = & F_i^n + A \left[ (F_{i+1}^{n+1} - 2F_i^{n+1} + F_{i-1}^{n+1}) + (F_{i+1}^n - 2F_i^n + F_{i-1}^n) \right] \\ & - B \left[ (F_{i+1}^{n+1} - F_{i-1}^{n+1}) + (F_{i+1}^n - F_{i-1}^n) \right] \end{aligned} \quad (5)$$

where  $A = \kappa \frac{\Delta t}{2\Delta z^2}$  and  $B = U \frac{\Delta t}{4\Delta z}$ .

## RESULTS AND DISCUSSION

We show the numerical results of the particle density at the interplanetary shock ( $z=0.5$  AU) at different times for a solar storm particle event (October 2000) in figure 2.



**Figure 2.** Results from numerical calculation at the shock at different times for Fe ions for the October 2000 event. Figure description.

We compare the results with measured spatial profiles of solar storm particles from the Ultra-Low-Energy Isotope Spectrometer (ULEIS) on board the Advanced Composition Explorer

(ACE) spacecraft. Our model previously provided a good fit to spectra of C, O, and Fe ions observed by ACE/ULEIS near the shock for 2 energetic storm particle (ESP) events in September 1999 and October 2000 [2].

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